

Investigation of the Depth Stability of the Furrow Former of a Saxaul Seed Planter

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Abstract

The article presents analytical expressions for determining the depth stability of the furrow former of a saxaul seed planter and identifies the factors that influence its performance.

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The **oval roller** is connected to the seeder frame through a **parallelogram mechanism**. Therefore, as shown in **Figure 1**, all the forces acting on it can be considered as applied to the lower movable hinge **O** of the parallelogram mechanism [1, pp. 118–120].

During operation, due to the **irregularities of the field surface** and the **variable physical and mechanical properties of the soil**, the normal force **N** acting on the working body and its horizontal (**N_x**) and vertical (**N_z**) components change continuously. As a result, the working body oscillates relative to the lower fixed hinge **O₁** of the parallelogram mechanism. This, in turn, leads to changes in the penetration depth of the working body into the soil and causes the furrow to form unevenly.

Therefore, the **amplitude of the oscillations** of the working body relative to the hinge **O₁** should be kept as small as possible. To address this problem, we derive and solve the **differential equation** describing the oscillations of the working body relative to the hinge **O₁**. For this purpose, the following assumptions are made [2, pp. 105–108]:

- The aggregate moves **in a straight line at a constant speed**;
- The **friction force** at the hinge **O₁** is small and has no significant influence on the oscillations of the working body;
- In its equilibrium state, the **lower longitudinal link** of the parallelogram mechanism is **horizontal** and, during operation, deviates from this position by only a **small angle**.

We take as the **generalized coordinate** the **deflection angle ε** of the lower longitudinal link of the parallelogram mechanism from the horizontal.

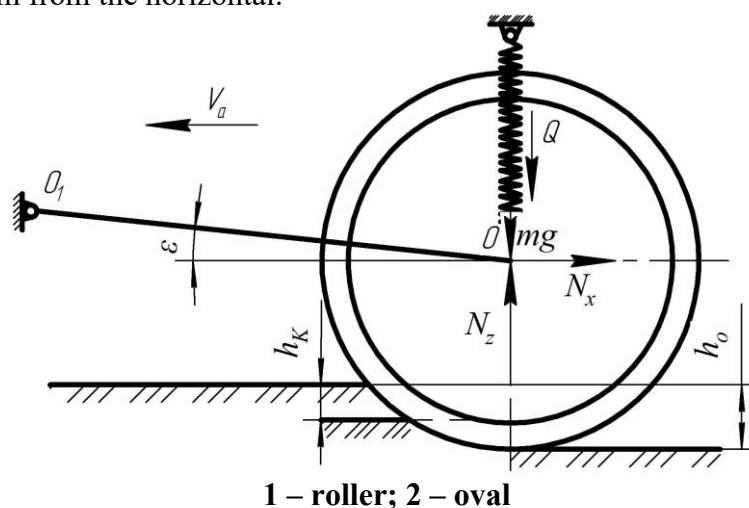


Figure 1. Diagram for studying the stability of the oval roller's penetration depth into the soil
 Using the equation of rotation of a rigid body about a fixed axis [3, pp. 279–280], we write:

$$J \frac{d^2 \varepsilon}{dt^2} = (N_z - Q - mg) l_n \cos \varepsilon - N_x l_n \sin \varepsilon, \quad (1)$$

Where

J – moment of inertia of the oval roller relative to the rotation axis O , $\text{kg} \cdot \text{m}^2$;

Q – spring pressure force, N ;

m – mass of the roller, kg ;

g – acceleration due to gravity, m/s^2 ;

l_n – length of the longitudinal link of the seeder's parallelogram mechanism, m .

Since the oscillation angle ε is small, we assume $\sin \varepsilon \approx \varepsilon$ and $\cos \varepsilon \approx 1$.

Taking this into account, equation (1) takes the following form:

$$J \frac{d^2 \varepsilon}{dt^2} = (N_z - Q - mg)l_n - N_x l_n \varepsilon. \quad (2)$$

We consider the soil reaction force N_z acting on the working body to consist of the **elastic component** N_y and the **resistance component** N_o , which depend on the **vertical displacement** and the **velocity** of this displacement, as well as the **variable force** N_t arising from **field surface irregularities** and the **changing physical and mechanical properties of the soil** [4, p. 230; 3, p. 58]. In this case —

$$N_z = N_y + N_o + N_t. \quad (3)$$

We substitute equation (3) into equation (2):

$$J \frac{d^2 \varepsilon}{dt^2} = (N_y + N_o + N_t - Q - mg)l_n - N_x l_n \varepsilon. \quad (4)$$

When the working body is in a **static equilibrium state**:

$$N_y = h_o C_n b; \quad (5)$$

$$N_o = 0; \quad (6)$$

$$Q = Q_0; \quad (7)$$

$$N_t = 0, \quad (8)$$

where

C_n – stiffness coefficient of the soil per unit working width of the implement, N/m^2 ;

$b = B_x$ – working width of the implement, m ;

Q_0 – initial compression force of the spring, N .

When the **parallelogram link** deviates from its equilibrium position by an angle ε —

$$N_y = (h_o - l_n \varepsilon) C_n b; \quad (9)$$

$$N_o = -b_n b l_n \frac{d\varepsilon}{dt}; \quad (10)$$

$$N_t = \Delta R_z(t) \quad (11)$$

and

$$Q = Q_0 + C_n l_n \varepsilon, \quad (12)$$

where

b_n – resistance coefficient of the soil per unit working width of the implement, $\text{N} \cdot \text{s/m}^2$;

C_n – stiffness coefficient of the spring.

If we substitute the values of N_y , N_o , N_t , and Q according to expressions (9)–(12) into equation (4), we obtain:

$$J \frac{d^2 \varepsilon}{dt^2} = \left[(h_o - l_n \varepsilon) C_n b - b_n b l_n \frac{d\varepsilon}{dt} + \Delta R_z(t) - (Q_0 + C_n l_n \varepsilon) - mg \right] l_n - N_x l_n \varepsilon. \quad (13)$$

When the working body is in **static equilibrium**:

$$(h_o C_n b - Q_0 - mg) l_n = 0. \quad (14)$$

Taking (14) into account, equation (13) takes the following form:

$$J \frac{d^2 \varepsilon}{dt^2} = \Delta R_z(t) l_n - C_n b l_n^2 \varepsilon - C_n l_n^2 \varepsilon - b_n b l_n^2 \frac{d\varepsilon}{dt} - N_x l_n \varepsilon \quad (15)$$

or

$$J \frac{d^2 \varepsilon}{dt^2} + b_n b l_n^2 \frac{d\varepsilon}{dt} + (N_x + C_n b l_n + C_n l_n) l_n \varepsilon = \Delta R_z(t) l_n. \quad (16)$$

Since the force N_x is variable, this equation is a **second-order differential equation with variable coefficients**.

From the **theory of vibrations**, it is known [5, pp. 110–112] that a system described by equation (16) may, in theory, exhibit **parametric vibrations**. However, because the soil has a **high damping capacity**, **parametric vibrations** of the working body are not observed. It mainly undergoes **forced oscillations** under the influence of the external force $\Delta R^z(t)$.

Taking this into account, we assume that the force R_x is **constant** and equal to its **average value**, and we consider the **forced oscillations** of the working body under the action of the force $\Delta R^z(t)$. Here we assume that the force $\Delta R^z(t)$ varies according to a **sinusoidal law**, that is —

$$\Delta R_z(t) = \Delta R \sin \omega t, \quad (17)$$

Where

ΔR – amplitude of the variable force, N;

ω – frequency of the variable force, 1/s.

Substituting equation (17) into equation (16), we write:

$$J \frac{d^2 \varepsilon}{dt^2} + b_n b l_n^2 \frac{d\varepsilon}{dt} + (N_x + C_n b l_n + C_n l_n) l_n \varepsilon = \Delta R l_n \sin \omega t \quad (18)$$

or

$$\frac{d^2 \varepsilon}{dt^2} + 2n \frac{d\varepsilon}{dt} + k^2 \varepsilon = H \sin \omega t, \quad (19)$$

where

$$n = \frac{b_n b l_n^2}{2J}; \quad k = \sqrt{\frac{(N_x + C_n b l_n + C_n l_n) l_n}{J}} \quad \text{ba} \quad H = \frac{\Delta R l_n}{J}.$$

It is known [6, p. 281] that the solution of equation (19), which describes the **forced oscillations of the working body**, is as follows:

$$\varepsilon(t) = \frac{H}{\sqrt{(k^2 - \omega^2)^2 + 4n^2 \omega^2}} \sin(\omega t - \Delta) \quad (20)$$

or, taking into account the adopted notations

$$\varepsilon(t) = \frac{\Delta R l_n \sin(\omega t - \Delta)}{J \sqrt{\left[\frac{(N_x + C_n b l_n + C_n l_n) l_n}{J} - \omega^2 \right]^2 + \left(\frac{b_n b l_n^2}{J} \right)^2} \omega^2}, \quad (21)$$

where

$$\Delta = \arctg \frac{b_n b l_n^2 \omega}{(N_x + C_n b l_n + C_n l_n) l_n - J \omega^2}.$$

The **maximum deflection angle** of the longitudinal link from its equilibrium position, according to equation (21), is equal to:

$$\varepsilon_{\max} = \frac{\Delta R l_n}{J \sqrt{\left[\frac{(N_x + C_n b l_n + C_n l_n) l_n}{J} - \omega^2 \right]^2 + \left(\frac{b_n b l_n^2}{J} \right)^2 \omega^2}}. \quad (22)$$

Expressions (21) and (22) show that the **stability of the working body in furrow formation**, and therefore the **quality of the furrow**, depends on the **moment of inertia of the roller**, the **length of the longitudinal link of the seeder**, the **spring pressure force of the parallelogram mechanism**, the **amplitude and frequency of the variable force**, as well as the **physical and mechanical properties of the soil**. Under the given working conditions and specific parameter values, the required stability of the working body in forming the furrow can mainly be ensured by the **proper selection of the spring pressure force** of the parallelogram mechanism.

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